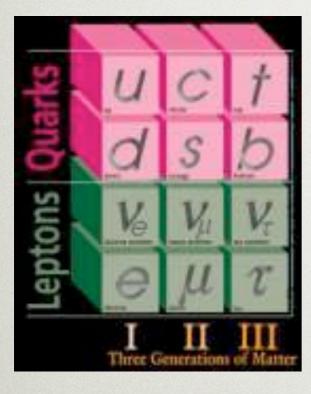
DIRECT DETECTION UPDATE AND CONNECTION TO INTENSITY EXPERIMENTS

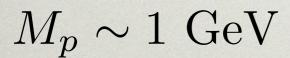
KATHRYN M. ZUREK
UNIVERSITY OF MICHIGAN

PARADIGM SHIFT

Our thinking has shifted



From a single, stable weakly interacting particle
(WIMP, axion)



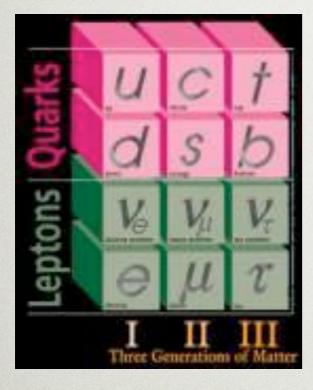
Standard Model

?

...to a hidden world with multiple states, new interactions

HIDDEN DARK WORLDS

Our thinking has shifted



From a single, stable weakly interacting particle
(WIMP, axion)

A

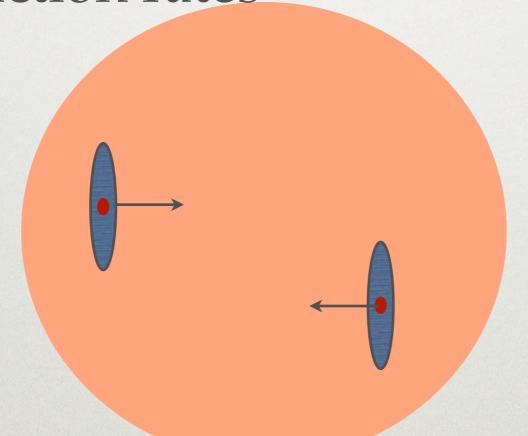
 $M_p \sim 1 \text{ GeV}$

Standard Model

...to a hidden world with multiple states, new interactions

WHY THE (SUB-)WEAK SCALE IS COMPELLING

Abundance of new stable states set by interaction rates

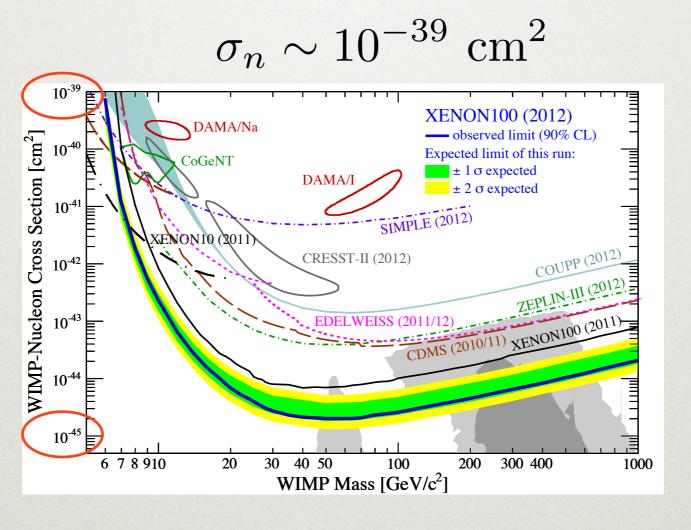


Freeze-out

$$\Gamma = n\sigma v = H$$
 $\rightarrow \sigma \sim rac{1}{ ext{few TeV}^2}$

SUB-WEAKLY INTERACTING MASSIVE PARTICLES

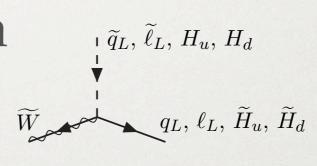
Scattering through the Z boson: ruled out

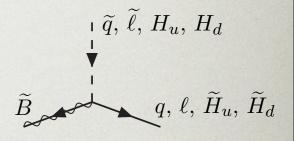


Next important benchmark: Scattering through the Higgs

$$\sigma_n \sim 10^{-45-46} \text{ cm}^2$$

- Make the Neutralino a pure state -- coupling \widetilde{W} $q_L, \ell_L, \widetilde{H}_u, \widetilde{H}_d$ \widetilde{B} $q, \ell, \widetilde{H}_u, \widetilde{H}_d$ to Higgs vanishes
- can be probed by indirect detection



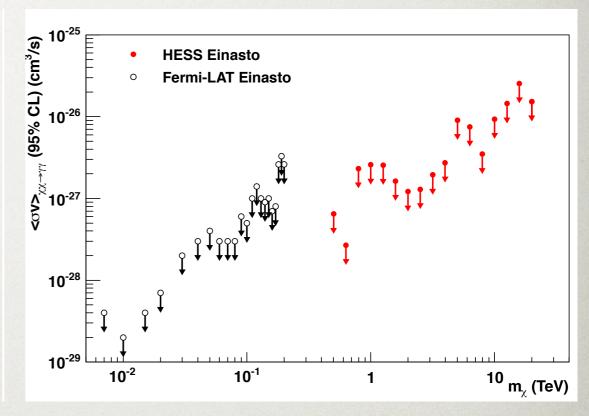


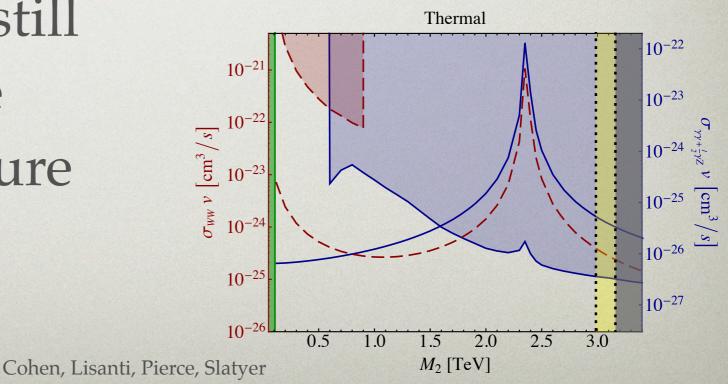
 However, Wino and Higgsino pure states

$$\chi$$
 χ
 χ
 χ
 χ
 χ
 χ

$$\langle \sigma v \rangle \sim \left(\frac{2 \text{ TeV}}{m_{\gamma}}\right)^2 10^{-26} \text{cm}^3/\text{s}$$

- Thermal Wino ruled out
- Thermal Higgsino still allowed, but can be ruled out in the future





- Bino escapes
- Pay a fine-tuning price

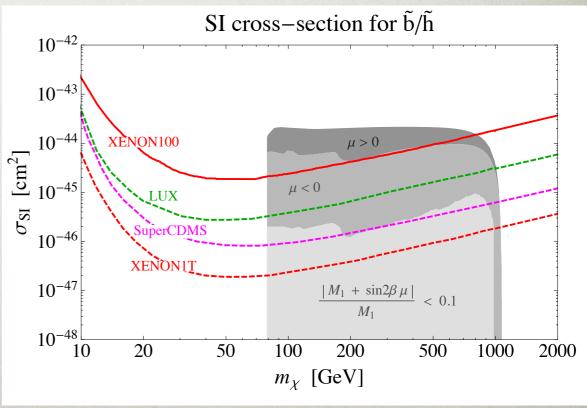
$$\mu \gg M_1 \sim m_{wk}$$

$$m_Z^2 = \frac{|m_{H_d}^2 - m_{H_u}^2|}{\sqrt{1 - \sin^2(2\beta)}} - m_{H_u}^2 - m_{H_d}^2 - 2|\mu|^2$$

- Tune away the coupling to the Higgs
- Smaller cross-sections
 correspond to more
 tuning in the neutralino
 components

\mathbf{m}_χ	condition
M_1	$M_1 + \mu \sin 2\beta = 0$
M_2	$M_2 + \mu \sin 2\beta = 0$
$-\mu$	$\tan \beta = 1$
M_2	$M_1 = M_2$

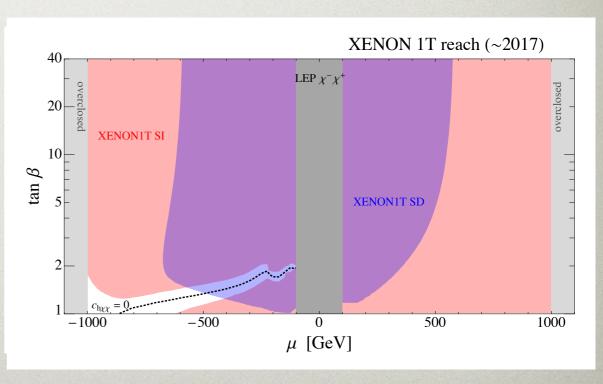
Cheung, Hall, Pinner, Ruderman



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Cheung, Hall, Pinner, Ruderman

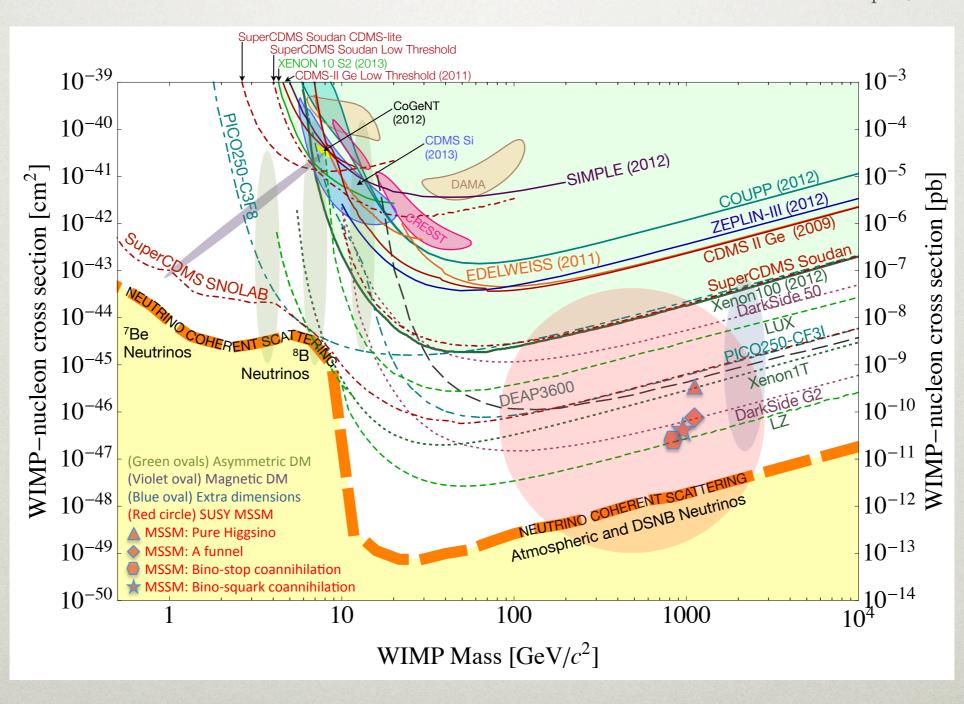


WHEN SHOULD WE START LOOKING ELSEWHERE?

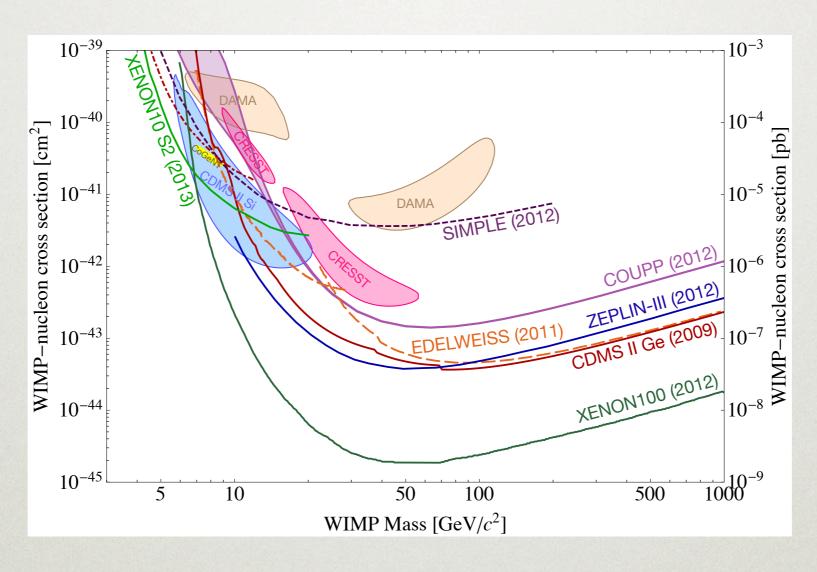
- Cannot kill neutralino DM, but paradigm does become increasingly tuned
- Somewhat below Higgs pole -- Neutrino background?
- Well-motivated candidates that are much less costly to probe
- Light WIMPs

TERRA INCOGNITA

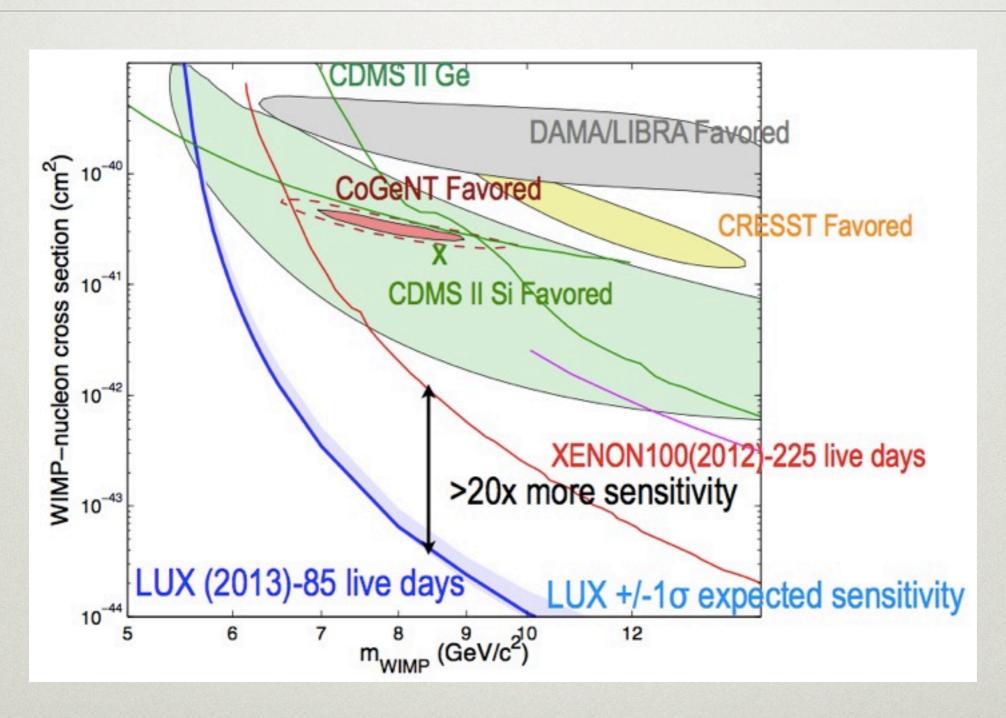
CF1 Snowmass report, 1310.8327



CURRENT SENSITIVITY LIMITED



ANOMALIES AND LUX



UNCERTAINTIES

- Experiment: Result assumes a particular choice of the energy calibration
- Theory: Also assumes spinindependent, momentum-independent scattering
- How do the results fare under more general assumptions?

OPERATOR UNCERTAINTIES

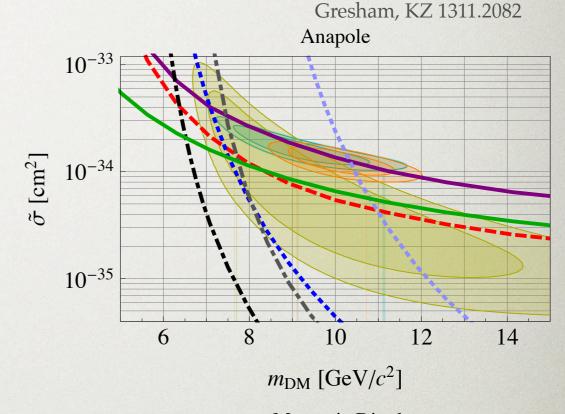
 Anapole and Dipole operators do best job, but neither escapes constraints

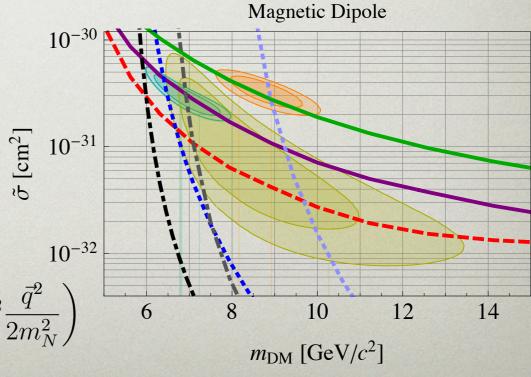
$$\mathcal{O}_a = \bar{\chi} \gamma^{\mu} \gamma_5 \chi A_{\mu}$$

$$\mathcal{O}_d = \bar{\chi} \sigma^{\mu\nu} \chi F_{\mu\nu} / \Lambda$$

$$\sigma_N^a = f_a^2 \frac{\mu_N^2}{\pi M^4} \left(Z^2 F^2(A; \vec{q}^2) \left(\vec{v}^2 - \frac{\vec{q}^2}{4\mu_N^2} \right) + \frac{J+1}{3J} g_N^2 A^2 \frac{\vec{q}^2}{2m_N^2} \right)$$

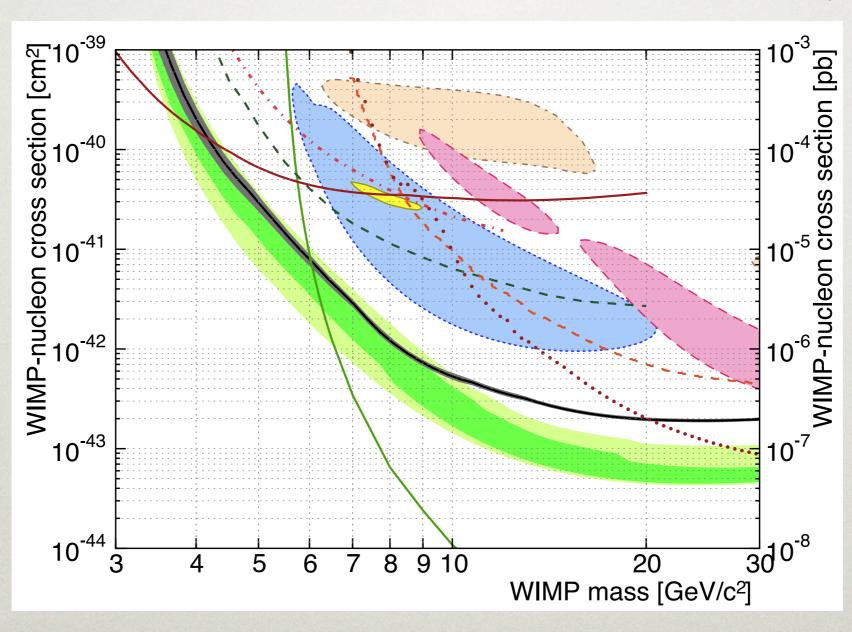
$$\sigma_N^d = f_d^2 \frac{\mu_N^2}{\pi M^4} \frac{\vec{q}^2}{\Lambda^2} \left(Z^2 F^2(A; \vec{q}^2) \left(\vec{v}^2 - \frac{\vec{q}^2}{4\mu_N^2} + \frac{\vec{q}^2}{4m_{\rm DM}^2} \right) + \frac{J+1}{3J} g_N^2 A^2 \frac{\vec{q}^2}{2m_N^2} \right)$$
10





SUPERCDMS: THE NEARLY FINAL WORD

1402.7137



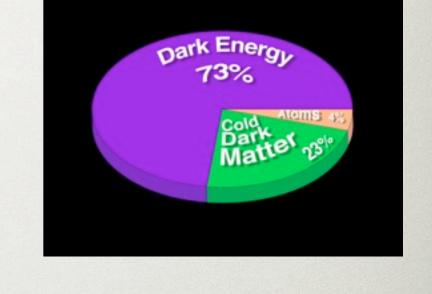
ANOMALIES PROBABLY NOT DUE TO DM

- But must be careful not to throw baby out with bath water
- Low mass DM is motivated theoretically, and does not necessarily predict excluded cross-sections



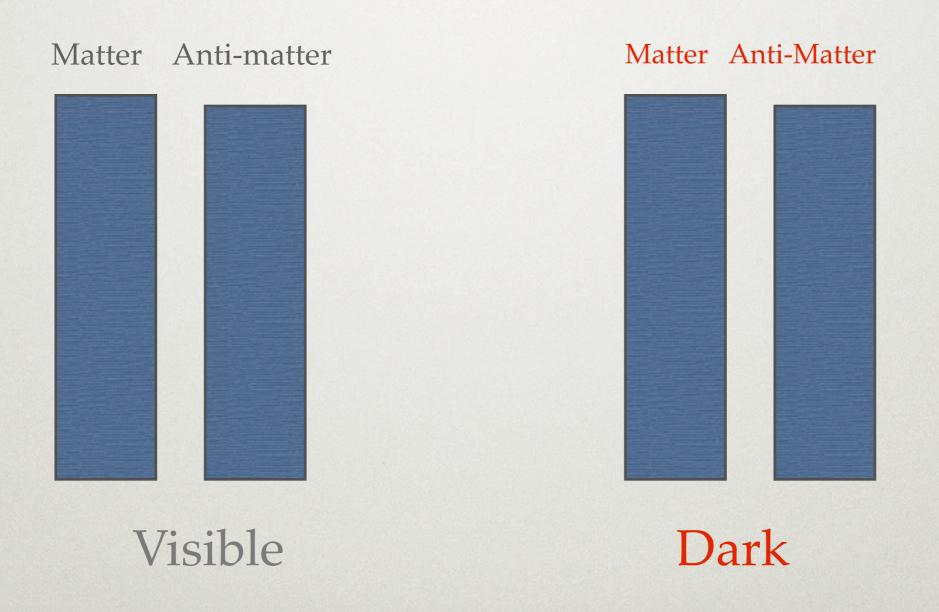
LIGHT WIMPS: ASYMMETRIC DARK MATTER

- Standard picture: freeze-out of annihilation; baryon and DM number unrelated
- Accidental, or dynamically related?



Experimentally, $\Omega_{DM} \approx 5\Omega_b$ Mechanism $n_{DM} \approx n_b$ $m_{DM} \approx 5m_p$

CHEMICAL POTENTIAL DARK MATTER



WHAT DOES AN ADM MODEL DO?

KZ, 1308.0338

- 1. Share an asymmetry between the visible and dark sectors
- 2. Decouple transfer mechanism to separately freeze-in the asymmetries in both sectors
- 3. Annihilate the symmetric abundance

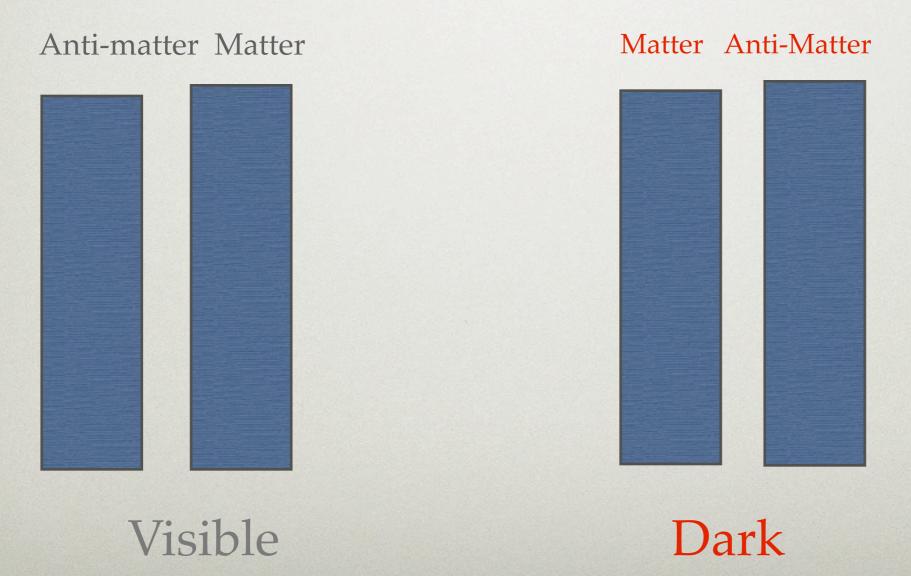
$$n_X - n_{\bar{X}} \sim n_b - n_{\bar{b}}$$



 $m_X \sim 5 m_p \simeq 5 \text{ GeV}$

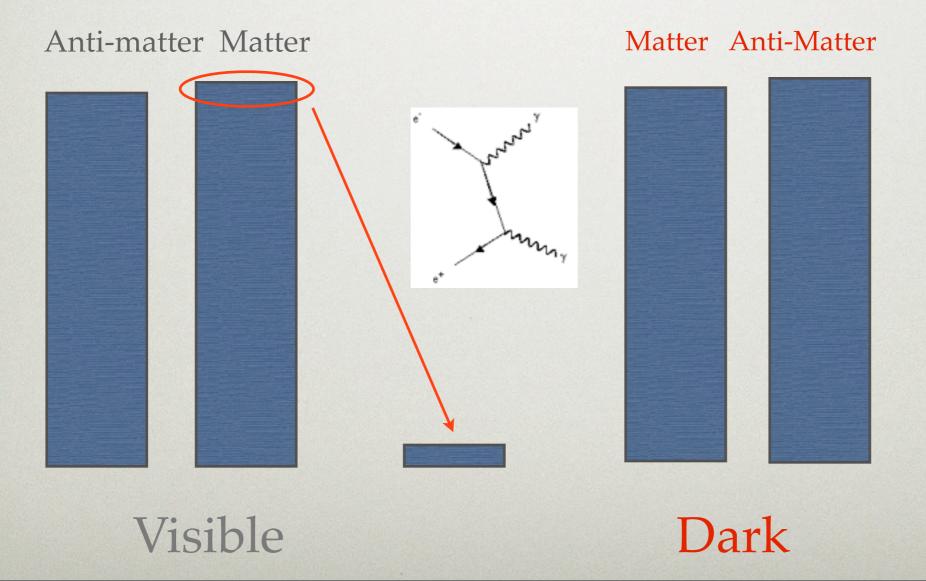
3. ANNIHILATING

• While it doesn't directly probe the asymmetry mechanism, it is more likely this physics is at a low scale which we can probe.



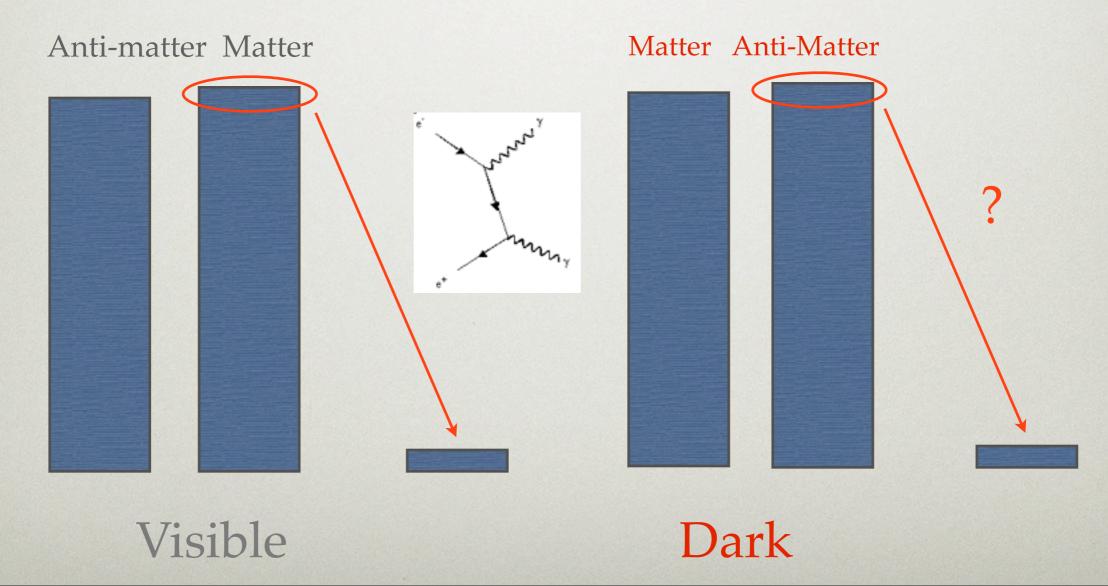
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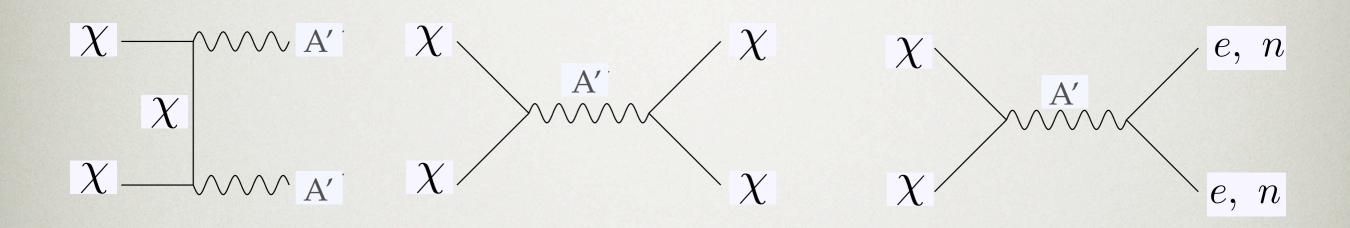


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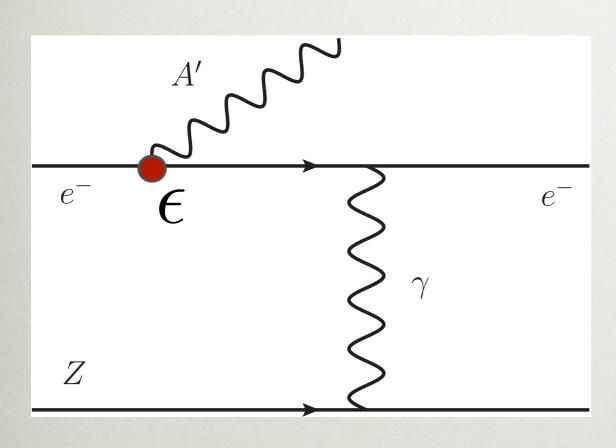


DARK FORCES AND DM INTERACTIONS

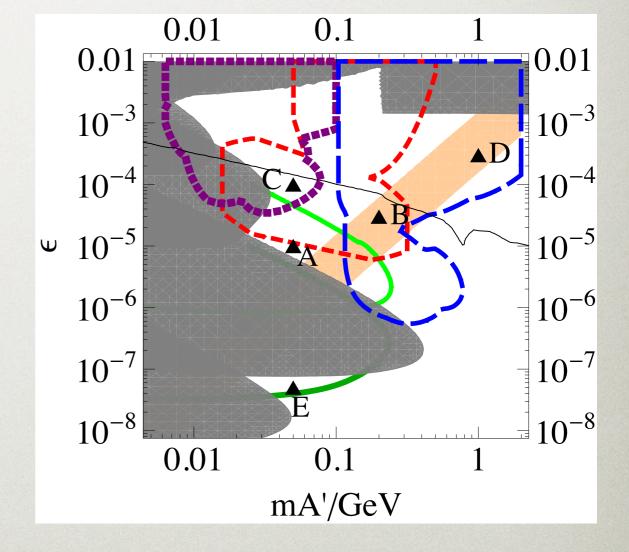


- Dark Forces Very Important for Light Dark Matter!
- May also be important for structure of DM halos
- May be important for DM direct detection and collider searches

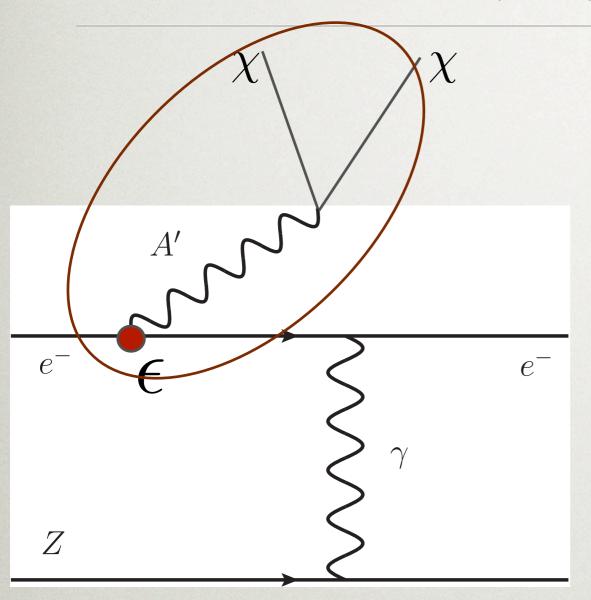
LOW ENERGY ACCELERATOR CONSTRAINTS

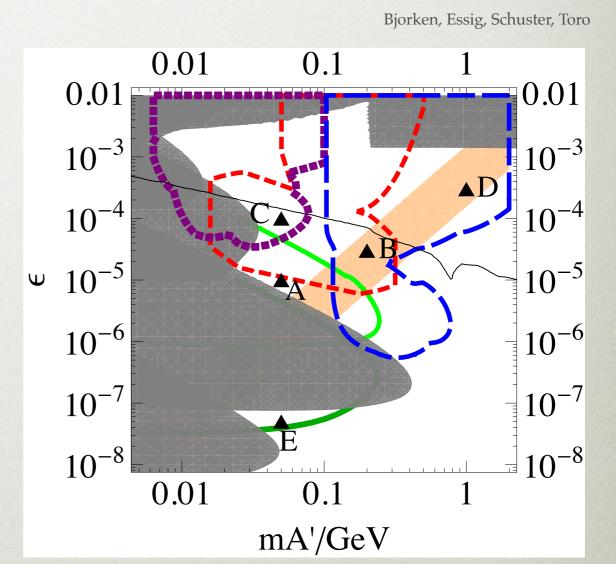


Bjorken, Essig, Schuster, Toro

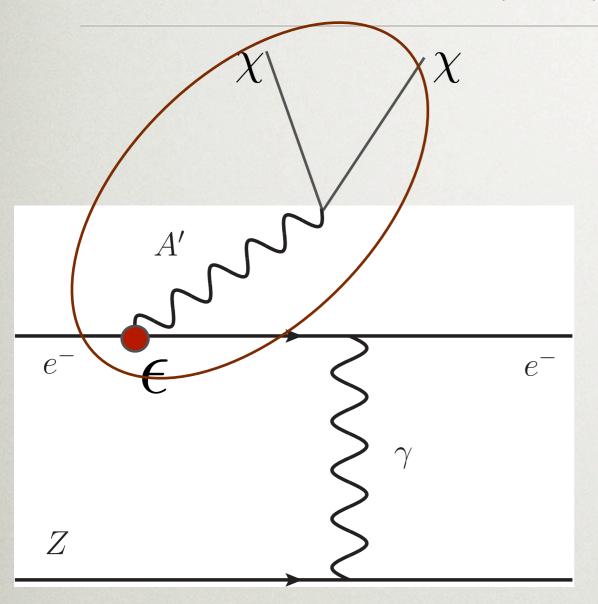


TRANSLATE TO DIRECT DETECTION

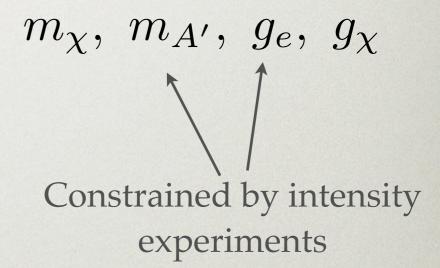




TRANSLATE TO DIRECT DETECTION

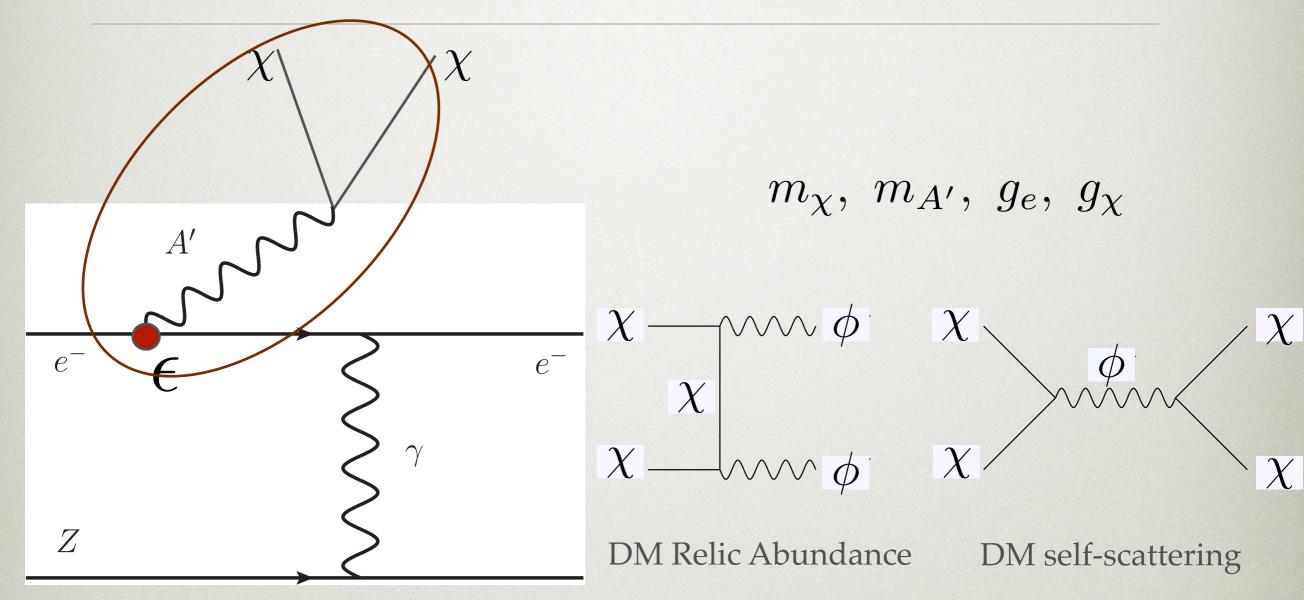


Ingredients:



Other complementary searches for other two parameters?

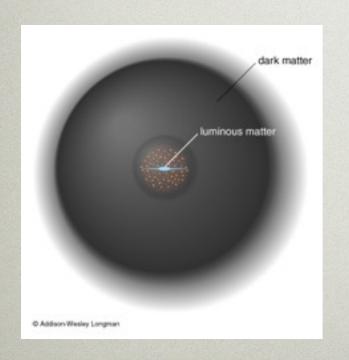
TRANSLATE TO DIRECT DETECTION



Can we connect dark photon searches to direct detection and other astrophysical observables?

DARK MATTER SELF-SCATTERING

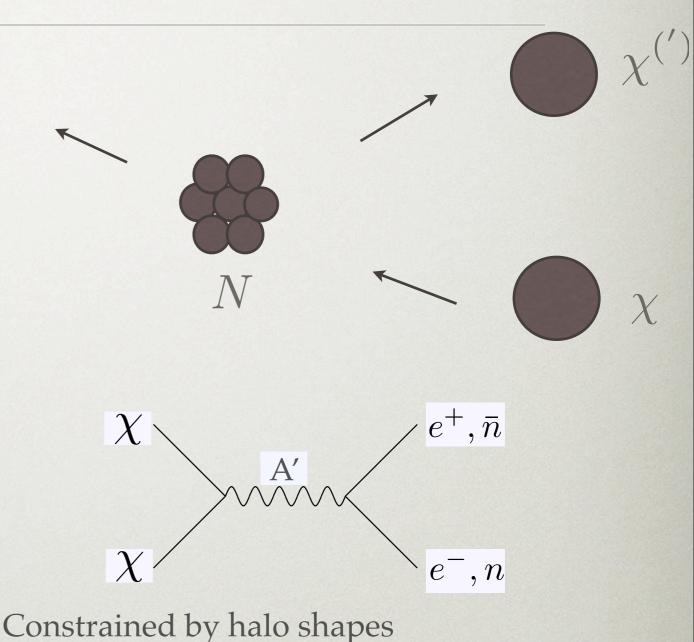
• Dark matter self-coupling changes the shape of a dark matter halo (such as the milky way halo) - we can extract constraints on coupling g_{χ}



$$\sigma_{\chi\chi} \approx \frac{g_{\chi}^4 m_{\chi}^2}{4\pi m_{A'}^4} \lesssim 4.4 \times 10^{-27} \text{ cm}^2 \left(\frac{m_{\chi}}{1 \text{ GeV}}\right)$$

CONNECTION TO DIRECT DETECTION

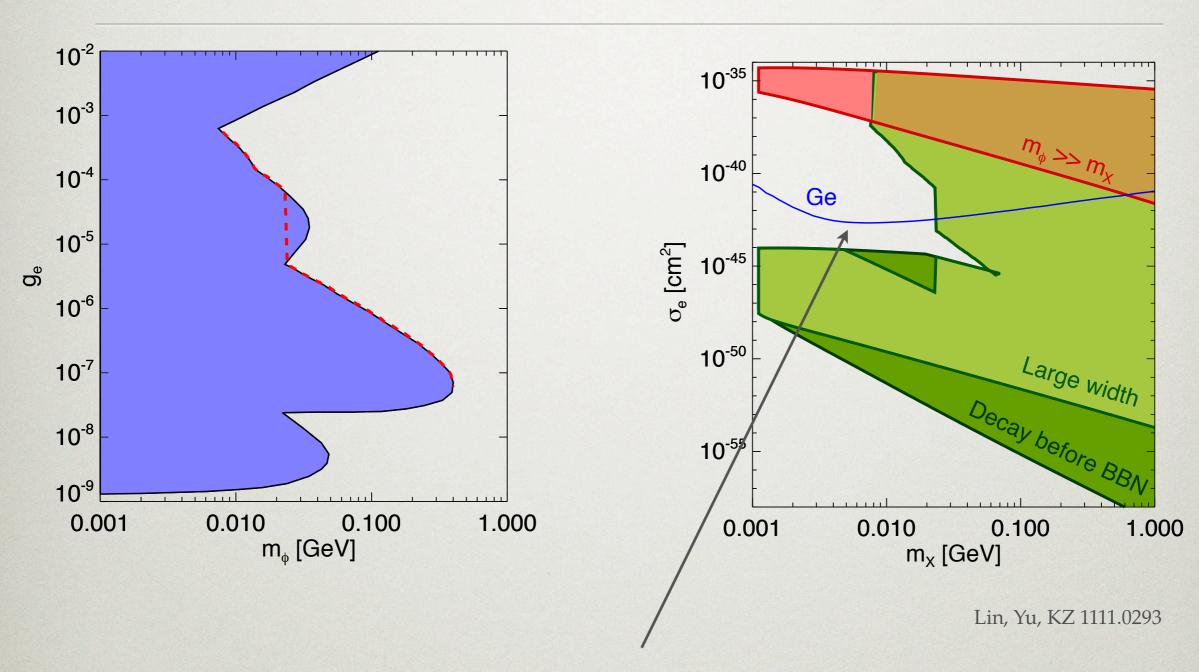
Can now take
 constraints from
 heavy photon
 searches + halo
 shapes to map to
 direct detection
 experiments



 $\sigma_n pprox rac{g_\chi^2 g_n^2 \mu_n^2}{\pi m_{A'}^4} \qquad \sigma_e pprox rac{g_\chi^2 g_e^2 \mu_e^2}{\pi m_{A'}^4}$

Constrained by heavy photon search

MAP INTO DIRECT DETECTION PLANE

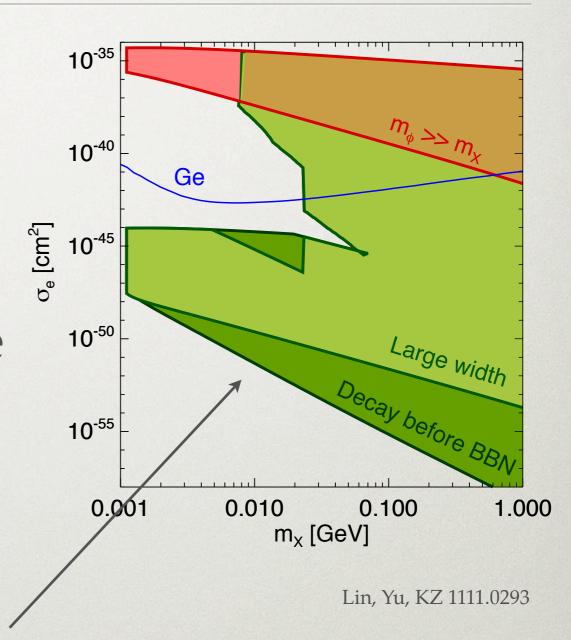


Projected maximum sensitivity of direct detection experiment

Cut-out gives combined constraints of beam dump + supernova + g-2

MAP INTO DIRECT DETECTION PLANE

Note that the lower bound of the theory parameter space is totally out of reach of any experiment! Can we do better?

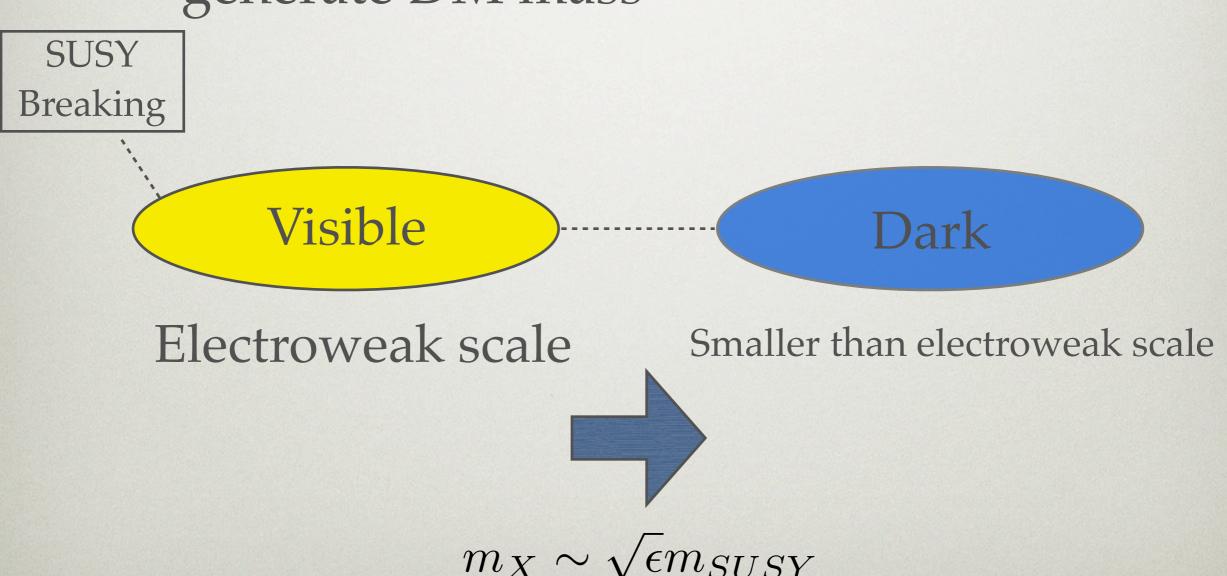


Require A' to decay before BBN

$$g_e \gtrsim 5 \times 10^{-11} \sqrt{10 \text{ MeV}/m_{A'}}$$

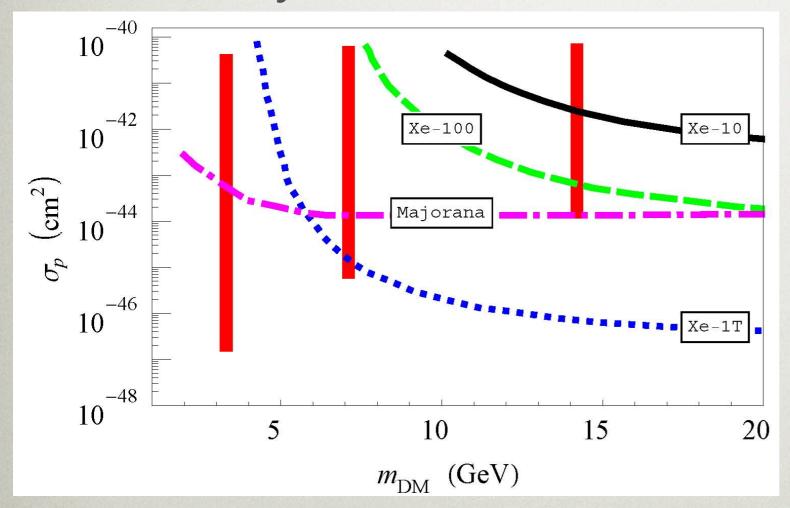
PARTICULAR MODELS CAN BE MUCH MORE PREDICTIVE!

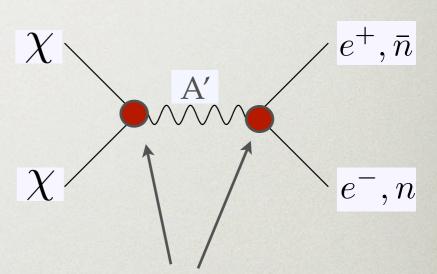
 Goal: explain why GeV? Dynamically generate DM mass



PARTICULAR MODELS CAN BE MUCH MORE PREDICTIVE!

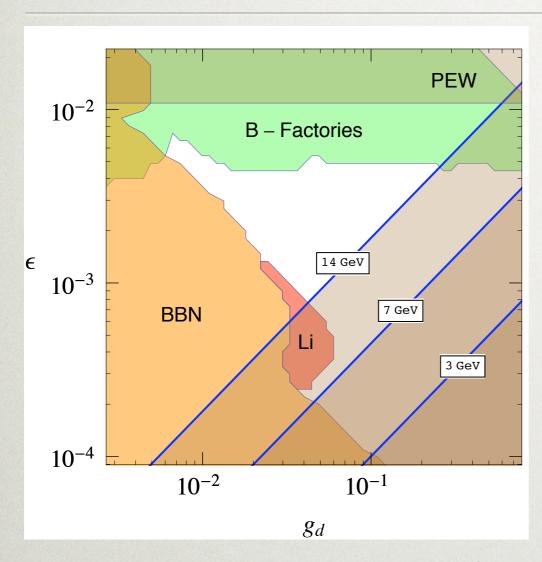
 Predict DD cross-section for Asymmetric Dark Matter!

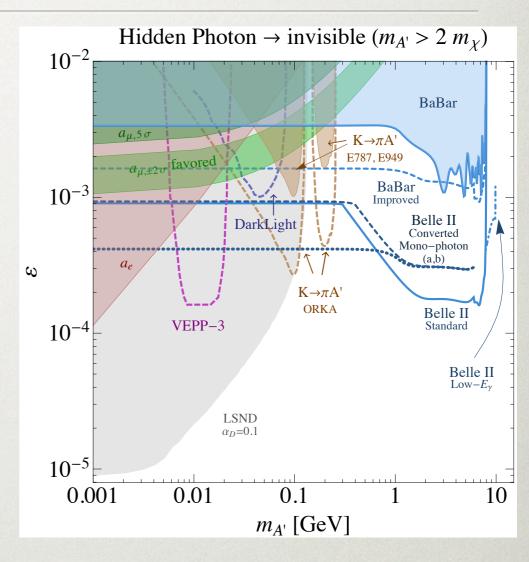




Coupling predicted by setting mass scale in DM sector!

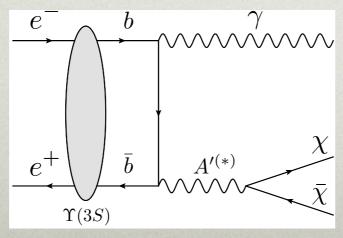
ALSO PROBED BY INTENSITY EXPERIMENTS

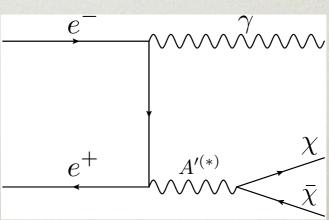




Cohen, Phalen, Pierce, KZ 1005.1655

Essig, Mardon, Papucci, Volansky, Zhong 1309.5084





SUMMARY

- In the last 7-10 years, particle theory has undergone a paradigm shift from sole focus on weak scale processes
- A key aspect of this paradigm shift is towards searching for light hidden sectors
- This light hidden sector may play a key role in the dynamics of the DM

SUMMARY

- Well-motivated models -- Asymmetric
 Dark Matter in particular
- Intensity experiments are complementary to direct detection and astrophysical probes
- Many probes coming online in next years